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A Technique to Ease the Fabrication Tolerance of Integrated Optical Power Splitters

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A Technique to Ease the Fabrication Tolerance of Integrated Optical Power Splitters

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Abstract

Fabrication tolerances and sidewall scattering losses in self-imaging waveguide devices are ameliorated by a partial-etch fabrication technique. Using a modal decomposition model, we find that the self-imaging plane's depth of focus increases with a reduction in etch depth. A broad depth of focus in the self-image plane relaxes the fabrication tolerance of the device's critical width dimension for a specified device performance. Trade-offs for this increased depth of focus include a modest increase in device length and a slight reduction in peak coupling efficiency.

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1. Introduction

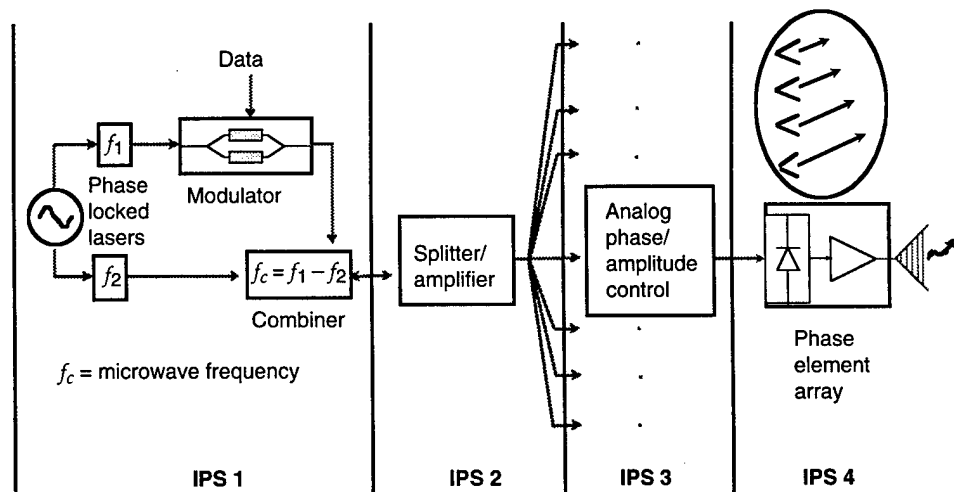
Research is under way at the Communications and Electronics Command (CECOM) and the Army Research Laboratory (ARL) in the development of integrated photonic subsystems (IPSs) for use in the optical control of phased-array antennas for communications on the move. Figure 1 shows a block diagram of the phased-array controller. In the figure, the module labeled IPS 2 provides the first level of signal splitting required to feed the antenna, which might vary from hundreds to thousands of elements. ARL is developing optical power-splitting devices to be used in IPS 2.

Desired characteristics of a $1 \times N$ splitter for this application include a value for N of about 16, the ability to split both the transverse electric (TE) and transverse magnetic (TM) modes, low crosstalk between the two polarization modes, uniformity among the split channels, low throughput loss, compact size, and ease of manufacture. For $N > 4$, fiber-optic polarization preserving splitters are impracticable. Therefore, the splitting must be performed on an integrated optical substrate. Numerous passive integrated optical beamsplitting techniques have been proposed and demonstrated. These include but are not limited to Y-junction branching waveguides, evanescent field directional couplers, computer-generated waveguide holograms, multimode interference (MMI) splitters, and radiative power splitters. For achieving the desired device characteristics, the MMI approach holds the greatest promise.

The practical implementation of MMI devices in guided-wave architectures largely depends on the device's fabrication tolerance, as defined by such performance metrics as excess optical loss. Using the paraxial approximation for strongly guided (i.e., deeply etched) structures, Besse et al [1] derived a closed-form approximation for the critical width dimension of the MMI region. For these deeply etched devices, they found that the fabrication tolerance is independent of the splitting ratio N and proportional to the output channel separation D .

A partial or shallow etch of the MMI device, however, is advantageous in many devices. For example, reduced sidewall interaction results in lower

Figure 1. Architecture for optical phased-array antenna control. Architecture is subdivided into four integrated photonic subsystems (IPSs).



excess loss and decreased nonradiative surface recombination in waveguide ring lasers [2]. Berry and Burke [3] used the discrete spectral index method to predict the self-imaging length and throughput of MMI devices as a function of etch depth. Shortly thereafter, 1×16 splitters with high throughput and good uniformity were demonstrated that were built by the partial-etch technique [4].

Since MMI devices are based on the principle of Talbot imaging (also known as self-imaging), the imaging plane's depth of focus can significantly affect the device fabrication tolerance. Recently, Smit et al [5] reported an increase in image plane focal depth with an increase in input rib width for deeply etched structures. In addition, excess loss was reduced, since a smaller fraction of the signal was contained in the higher order modes.

In this report, we review the results of a theoretical investigation of the depth of focus dependence on etch depth in MMI devices. We find that a shallow etch depth yields an extended depth of focus and thus a broader fabrication tolerance, at the expense of a slight penalty in throughput. This throughput loss results from a reduction in the number of modes supported within the MMI region, and we show that this loss is negligible for etch depths beyond mode cutoff.

2. Theory

Since a 1×1 MMI device produces a single self-image of the input, this configuration (fig. 2) is the simplest for investigating the effect of etch depth on imaging plane depth of focus. We make the following assumptions concerning the MMI structure under investigation (fig. 2): (1) the depth of focus is analyzed at the first single self-image plane, (2) the input rib, MMI region, and output rib are all defined in a single etch step, and (3) the device sidewalls are vertical.

Figure 3 shows the refractive index and thickness values used to define the transverse waveguide structure. These values correspond to an InGaAs/InAlAs waveguide operating at $1.319 \mu\text{m}$. Although these parameters vary from one waveguide structure to the next, our analysis is presented in terms of mode index difference between the MMI region and the surrounding etched regions; this approach allows us to generalize to other structures.

The etch depth of those areas surrounding the MMI region determines the mode index difference. Together with the MMI region width, this index difference determines the number of lateral modes supported within the MMI region. These modes will be excited to varying degrees by an input to the MMI region. Since each mode propagates with a slightly different phase velocity, the modes become dephased. At the place where the accumulated phased differences among the modes reach an integral multiple of 2π , a self-image of the input to the MMI region is reconstructed; i.e., the image is created solely by virtue of diffraction. No optical elements are required.

To reduce the three-dimensional structure of figure 2 to a two-dimensional structure, we have modeled the self-image formation by

Figure 2. Perspective view of a 1×1 multimode interference device.

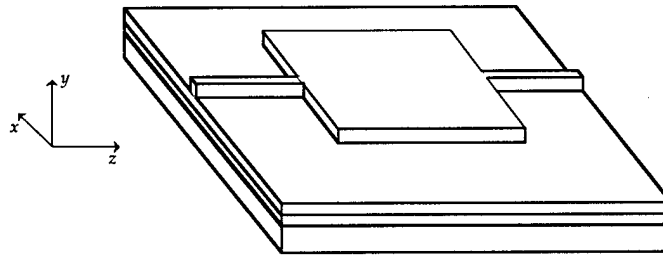
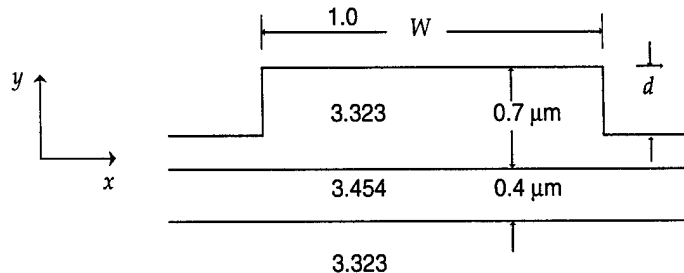


Figure 3. Transverse waveguide structure with film indices and thickness at $1.319 \mu\text{m}$: cross section through MMI region.



using the effective index method along the transverse waveguide dimension. The lateral modes of the MMI region are then calculated. The rib waveguide mode is decomposed into these lateral MMI modes (which are propagated the length of the MMI region), and the self-image is constructed [6,7]. As the etch depth is increased, a greater number of lateral modes are supported, not only in the MMI region, but also in the input rib waveguide. However, in our calculations, we assume that only the fundamental mode of the input rib waveguide is excited.

3. Results

The gray-scale contour plots of the MMI region's electric field amplitude (fig. 4) show the field evolution through the MMI region for three characteristic etch depths: very shallow, at mode cutoff, and very deep. The input rib width is fixed at $2\text{ }\mu\text{m}$, and the MMI region width is fixed at $10\text{ }\mu\text{m}$ for each of the etch depths. Two points are immediately apparent from figure 4. First, the MMI region's self-imaging length is longer for shallow etch depths; second, the depth of focus and input field width are greater with shallower etch depths. These observations of a waveguide self-imaging system parallel the general behavior of a single-lens imaging system. For a fixed aperture size, as the imaging distance is increased, both the depth of focus and the focused spot size are increased [8].

Berry and Burke [3] used the discrete spectral index method to predict the position of the self-imaging plane as a function of etch depth. Our results, based on the effective index method and modal propagation, corroborate their findings. The increase in self-imaging plane length with decreasing etch depth can be attributed to the *effective* width of the MMI region. The

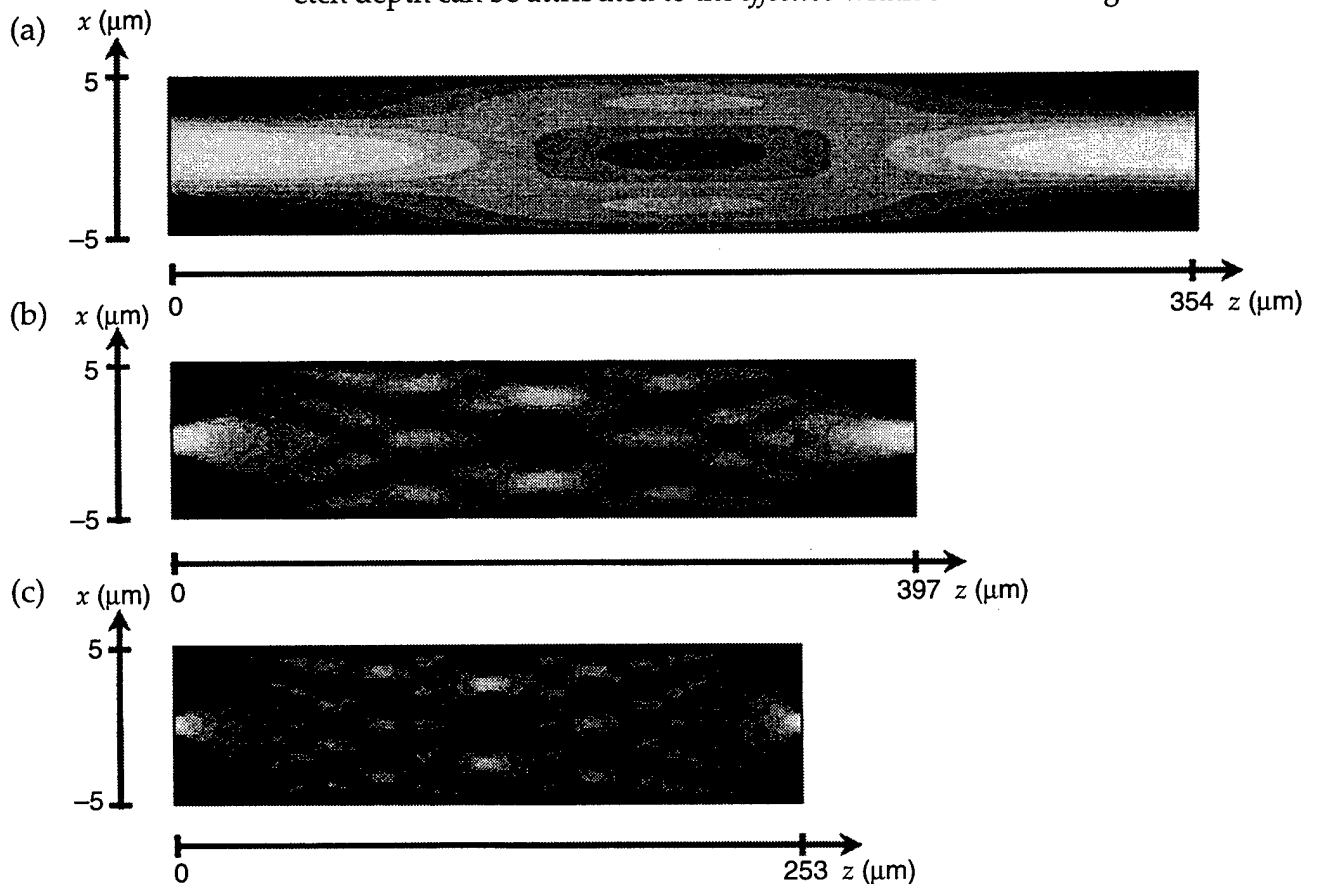


Figure 4. Top view of MMI region showing a gray-scale plot of TE field. (a) Etch depth $0.52\text{ }\mu\text{m}$; only 3 lateral modes supported in MMI region. (b) Etch depth $0.82\text{ }\mu\text{m}$ (transverse mode is at cutoff); 10 lateral modes supported in MMI region. (c) Semi-infinite etch depth assumed (i.e., lateral mode index surrounding MMI region and input/output rib waveguides is unity); 49 lateral modes supported in MMI region.

effective width W_e is equal to the MMI region's physical width W , corrected by the Goos-Hänchen penetration depth [9]:

$$W_e = W + \left(\frac{\lambda}{\pi}\right) \left(\frac{n_{lat}}{n_{mmi}}\right)^{2\sigma} \frac{1}{\sqrt{n_{mmi}^2 - n_{lat}^2}}, \quad (1)$$

where λ is the free space wavelength, n_{lat} and n_{mmi} are the transverse mode indices of the lateral (etched) and MMI (unetched) regions, respectively, $\sigma = 0$ for TE polarization, and $\sigma = 1$ for TM polarization. In this expression, we assigned a fixed effective width for all the lateral modes in the MMI region. (Our actual analysis is more rigorous.) Given this etch-depth-dependent effective width, the self-image plane positions of a $1 \times N$ center-fed MMI coupler are approximated by [10]

$$L \cong \frac{n_{mmi} W_e^2}{N\lambda}. \quad (2)$$

The analytical approximations of equations (1) and (2) are in good agreement with the more exact modal propagation analysis of the self-imaging length's etch-depth dependence (see fig. 5). As the MMI region width is increased, however, the analytical approximation breaks down.

The efficacy of MMI devices in photonic switching systems is contingent on the relative ease in fabricating high-throughput devices. The imaging plane depth of focus determines the fabrication tolerance on the MMI region's critical width dimension. In this report, we use the coupling efficiency from the self-image formed at the end of the MMI region into the output rib as the defining metric for depth of focus. We calculate the coupling efficiency by performing an overlap integral of the MMI field with the mode supported by the output rib. This efficiency is converted to excess device loss.

In figure 6, we plot the depth of focus (assuming a maximum permissible excess loss of 1 dB) versus the mode index difference between the MMI and lateral regions for TE polarization. For mode index differences less than that corresponding to mode cutoff in the lateral region, the depth of focus varies rapidly with mode index difference.

In figure 7, we used the effective index method to plot the depth of focus versus etch depth for depths down to mode cutoff. In the limit of a very deep etch, the effective index of the lateral region approaches unity, yielding an asymptotic limit for the depth of focus. We can approximate the MMI device fabrication tolerance by differentiating equation (2) with respect to the physical device width W and rearranging:

$$\Delta W = \frac{N\lambda \Delta L}{2n_{mmi} W_e}. \quad (3)$$

For our test structure, a deep etch yields a depth of focus of about $12 \mu\text{m}$, which requires a tolerance of $\pm 0.12 \mu\text{m}$ in the MMI region width. At mode cutoff, the depth of focus is increased to about $28 \mu\text{m}$, for a more easily manufacturable MMI width tolerance of $\pm 0.26 \mu\text{m}$.

Figure 5. Comparison of 1×1 self-imaging lengths as a function of etch depth for modal propagation and approximate analytical solutions.

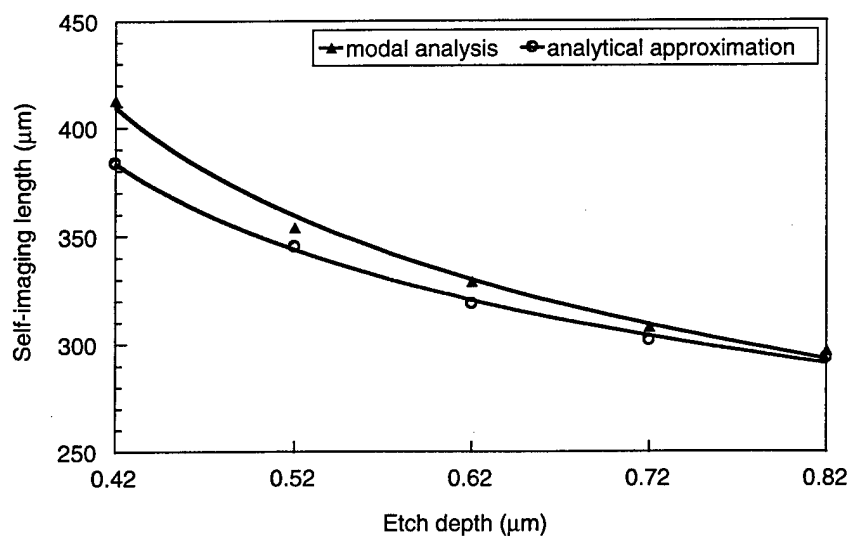


Figure 6. Depth of focus versus mode index difference between MMI region and surrounding etched regions. Depth of focus is determined at an excess coupling loss of 1 dB.

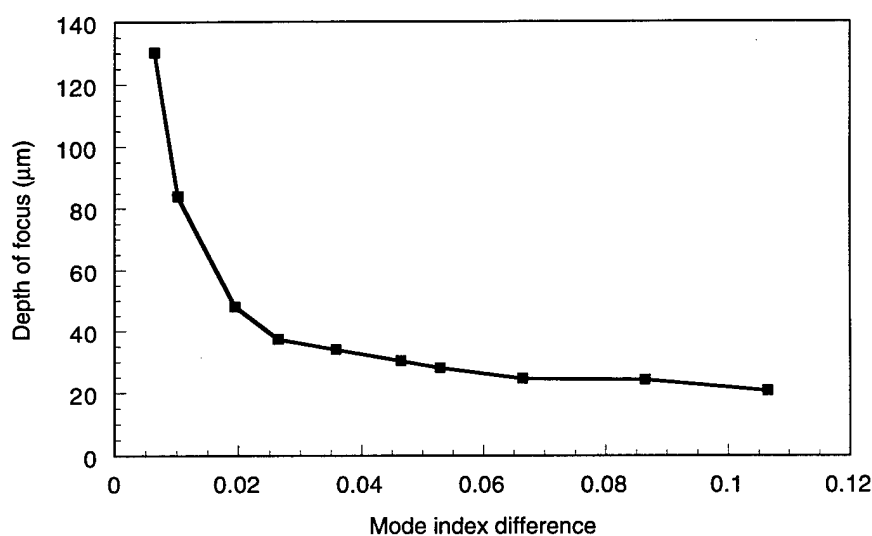
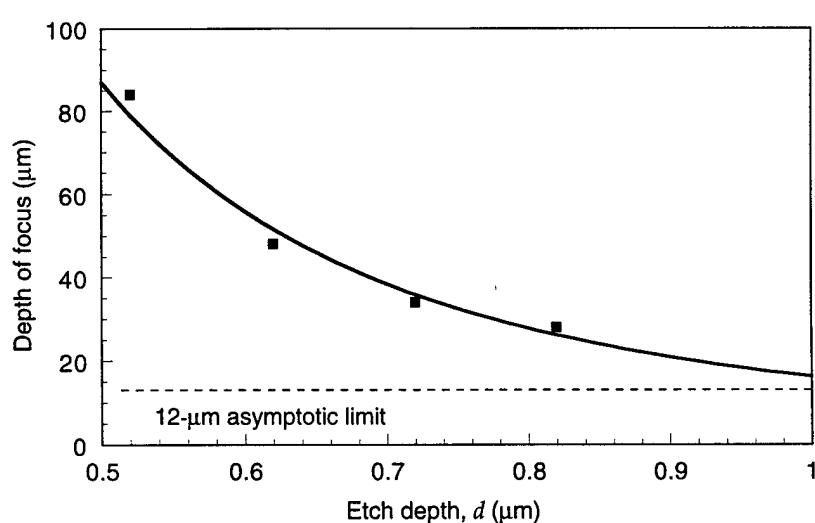


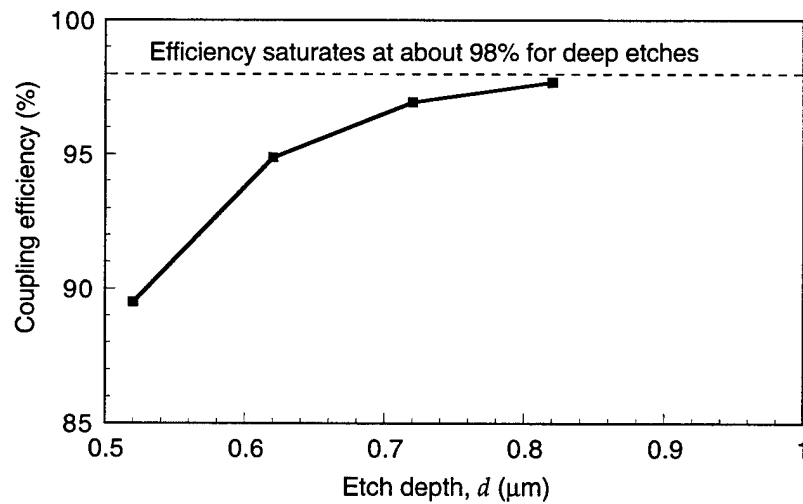
Figure 7. Depth of focus versus etch depth obtained by effective index method. Depth of focus is determined at an excess coupling loss of 1 dB. For a semi-infinite etch depth, depth of focus asymptotically approaches 12 μm.



Although the data presented here correspond to the specific structure selected, the basic trend of an increased depth of focus (and thus fabrication tolerance with a decreased etch depth) is generalizable to arbitrary waveguide structures. As shown in figure 4, the shallower etch depths produce a broader fundamental mode in the fixed 2- μm -wide input rib. This extended depth of focus with a laterally wider input field is consistent with the findings of Smit et al [5], who investigated extending the depth of focus for deeply etched MMI devices by increasing the physical width of the input rib.

Several drawbacks arise from a reduced MMI device etch depth. We have already characterized the increase in device length due to the broader effective width as given by equations (1) and (2). Next, the etch depth affects self-image quality. With a smaller mode index difference between the MMI and laterally etched regions, fewer lateral modes are available for reconstructing the input for self-image formation. The consequence of the distorted image is to reduce the coupling efficiency into the output rib. However, figure 8 shows that a reduction in peak coupling efficiency from the distorted image into the output rib is negligible for etch depths to mode cutoff or greater. Finally, a shallower etch depth in the N output ribs requires a larger output rib separation to avoid mutual coupling among the ribs. If the output rib separation is accomplished with waveguide S-bends, a larger radius of curvature is required, because of the weaker mode confinement.

Figure 8. Maximum coupling efficiency of self-image into output rib waveguide as a function of etch depth. Transverse mode is cut off at etch depth of 0.82 μm .



4. Conclusions

The etch depth in MMI device fabrication is a design variable that can be exploited to improve device manufacturability. In the 1×1 test structure, our model predicted a doubling of the MMI width fabrication tolerance when the MMI device is etched to a depth where the transverse waveguide mode is just cut off, as opposed to the deep etches that are typically used. This doubling is accomplished through an increase in the image plane depth of focus. The concomitant distortion of the self-image is negligible for etches to mode cutoff or deeper. These advantages to varying the MMI device etch depth are traded off with an increase in device length and weaker confinement in the modes of the output rib array.

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